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A transient absorption study of allophycocyanin

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Abstract. Transient dynamics of allophycocyanin trimers and monomers are observed by using the pump-probe, transient absorption technique. The origin of spectral components of the transient absorption spectra is discussed in terms of both kinetics and spectroscopy. We find that the energy gap between the ground and excited states of the unexcited subunit of allophycocyanin monomer decreases via an interaction with another excited subunit. For allophycocyanin trimer, we find that the fast dynamics results from the fast internal conversion and the first excited state is the only one electronic state which can trap the final population.

Keywords. Allophycocyanin; transient absorption; induced absorption; stimulated emission; ground state bleaching; internal conversion.

1. Introduction

The phycobiliproteins, allophycocyanin (APC) and C-phycocyanin (C-PC), are presumably considered to play an important role in the energy transfer process of plants and algae. These proteins are open-chain tetrapyrrole-containing proteins located at the core of the phycobilisome in cyanobacteria.¹ The X-ray crystal structure of the trimer aggregation state exhibits C₃-symmetry. In particular, the structures of such proteins show a ring-like homotrimer of three monomers consisting of **a** and **b** polypeptide monomers as has been already discussed in references.^{2–4}

In this work, we shall obtain basic information on transient dynamics of APC trimers and monomers based on pump-probe, transient absorption measurements. An origin of spectral components of the transient absorption spectra will be discussed.

2. Experimental setup

The transient absorption (TA) experimental arrangement is shown in figure 1, and the ultrashort laser system is the same as used in our previous work.⁵ The pump energy is provided by the output of an OPA laser and attenuated to 30 nJ per pulse with a pulse duration time of 100 femtoseconds (fs) with a FWHM of 6~7 nm at a repetition rate of 20 kHz; while the probe beam is a continuum white light centered at 810 nm (460~

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1200 nm, ~1 nJ), which is generated by focusing a 810 nm laser light on a sapphire crystal. In front of the sample cell, the continuum white light is divided into two beams with equal energy; one is for a probe beam and another one is for a reference beam. The pump and probe beams are focused on the same spot of the sample in a quartz flow cell of 1 mm thickness and the reference beam is also focused on a different location in the flow cell. After the flow cell, the probe and reference beams are sent to a detector (Oriel MS257). Both probe and reference signals coming out from the monochromator and photomultiplier are directly sent to a personal computer for processing in the real time. After calibrating the absorbance intensity of probe beam by the reference beams, we obtain the difference absorbance (ΔA). In this case, ΔA is given as the absorbance difference between with and without pump beam. In the experiment, the APC, used herein, are isolated from cyanobacterium cultures of the *Anabaean variabilis*. The preparations and characterization methods for APC trimers and monomers are described in reference.⁶



Figure 1. The experimental setup of transient absorption measurements. M: minor, L: lens, P: prism, BS: beam splitter, TS: transition stage, RR: retro reflector, ST: stopper, ND: ND filter, CF: short cut off filter between 750 and 950 nm, PB: probe beam, RB: reference beam.

3. Results and discussions

In the pump-probe system, generally, at least three types of optical processes exist (1) ground state bleaching (GB), (2) stimulated emission (SE), and (3) induced absorption (IA). The total signal of the difference absorbance is given by the summation of the three signals. In our transient absorption measurements, the IA signal is positive (+), while the GB and SE signals are negative (–). Figure 2 shows three-dimensional, transient absorption spectra of APC monomer (figure 2a) and APC trimer (figure 2b). Here APC trimers are excited at 580 nm with parallel polarization, while APC monomers are excited at 570 nm with magic angel (54.7°) polarization. Note that in figure 2a worse S/N ratios at the pump wavelength (~570 nm) can be seen, which is due to scattered light of the pump pulse. Figure 2a shows a very slow decay component due to the population dynamics of the excited state during the observation time ~150 ps. The strong IA bands at both the red and blue sides of the TA spectra appear immediately after the pump, and these bands almost persist without decreasing their intensities during the observed delay time. While



Figure 2. The observed three-dimensional, transient absorption spectra of APC monomer (a) and trimer (b).

in figure 2b, the intensities of IA bands are found to decay with different time constants. Note that the intensity of IA bands depends on the pump power.

3.1 TA spectra in frequency domain

Figure 3 presents the observed transient absorption spectra of APC monomer (figure 3a) and trimer (figure 3b). The observed spectra are normalised at the maximum intensity in the negative intensity region and are shown as the function of probe wavelengths at several delay times. For comparison, the normalized absorption (dot line) and fluore-scence (dash line) spectra are also shown in figure 3. During 150 ps, figure 3a exhibits a merged band of both the ground state bleaching and stimulated emission (GB/SE) located between the absorption and fluorescence spectral regions. From figure 3a, one can see that the TA spectra are all the same during the delay time. On the other hand, the transient spectra of **a**subunit⁷ excited at 620 nm are very similar to those of APC monomer (pumped at 570 nm). The similarities indicate that the observed features of APC monomer simply result from individual **a** and **b**subunits and these two chromophores weakly interact with each other.



Figure 3. The normalized transient spectra at various delay times for APC monomer (a) and trimer (b).

It is interesting to compare the monomer spectra with that of the APC trimer. From figure 3b, we notice that (1) the TA spectra within the 640-670 nm region provides the merged spectra of ground state bleaching and stimulated emission (GB/SE); (2) the 640-625 nm and 660-680 nm regions exhibit induced absorption (IA) bands; (3) the 580-620 nm region shows the GB spectra; and (4) the 700-720 nm region provides SE. It is found that the GB/SE merged band appears within 2 ps and it is located at the steady-state fluorescence region. This band exhibits a red shift with an increase of the delay time.

3.2 The modeling of APC

To construct a microscopic model for the dynamics appearing in the observed time profiles of APC trimer and monomer, we attempt to extract IA spectra components by assuming that (1) GB and SE spectra components are directly associated with the steady-state absorption and fluorescence spectra and GB spectra components are independent of the delay time and (2) the blue and red region of TA spectra only contain GB and SE components.

3.2a APC monomer: We consider a microscopic model for APC monomer by assuming that the excitation takes place on **a**subunit or **b**subunit is given by $|\mathbf{a}_g\rangle|\mathbf{b}_g\rangle^{-\frac{\text{pump}}{\text{pump}}}|\mathbf{a}_g\rangle\mathbf{b}_e\rangle$ or $|\mathbf{a}_e\rangle|\mathbf{b}_g\rangle$. Note that two chromophores in the monomer are coupled weakly with each other. In our model, the electronic energies of $|\mathbf{a}_g\rangle|\mathbf{b}_e\rangle$ and $|\mathbf{a}_e\rangle|\mathbf{b}_g\rangle$ are the same and only one subunit is excited. For example, excitation of **a**subunit can lead to an interaction with the unexcited **b**subunit, modifying the energy structure of the **b**subunit. In this case, the energy gap between $|\mathbf{a}_e\rangle|\mathbf{b}_g\rangle$ and $|\mathbf{a}_e\rangle|\mathbf{b}_e\rangle$ (the second excited state in IA band) can be smaller. In other words, the net absorbance of TA spectra will show positive IA bands at the red side of the GB/SE band. This is due to the fact that the absorption spectra of chromophore in **b**subunit is slightly shifted after the **a**subunit being excited.

Using two fitting rate constants, $1/25 \text{ ps}^{-1}$ and $1/300 \text{ ps}^{-1}$, as the initial parameters, we can describe the vibration relaxation rate constant k_r of $|\mathbf{n}\rangle$ and the electronic relaxation rate constant k_e of $|\mathbf{a}_k\rangle|\mathbf{b}_k\rangle$, respectively. The population in each state is redefined as \mathbf{r}_1 for the hot vibrational state $|\mathbf{n}\rangle$ of $|\mathbf{a}\rangle|\mathbf{b}\rangle$, \mathbf{r}_2 for the bottom of electronic state $|\mathbf{a}_{\mathbf{k}}\rangle|\mathbf{b}_{\mathbf{k}}\rangle$, and \mathbf{r}_{0} for the ground state $|\mathbf{a}_{\mathbf{k}}\rangle|\mathbf{b}_{\mathbf{k}}\rangle$. Notice that $d\mathbf{r}_{1}/dt = -k_{r}\mathbf{r}_{1}$ and $d\mathbf{r}_2/dt = k_r \mathbf{r}_1 - k_e \mathbf{r}_2$. The absorbance A consisting of the contributions from \mathbf{r}_1 , \mathbf{r}_2 and \mathbf{r}_0 can be written as $A = (\mathbf{e}_{gb} + \mathbf{e}'_{gb})\mathbf{r}_0 + (\mathbf{e}'_{ia} - \mathbf{e}'_{se})\mathbf{r}_1 + (\mathbf{e}_{ia} - \mathbf{e}_{se})\mathbf{r}_2$, where $\mathbf{r} = \mathbf{r}_1 + \mathbf{r}_2 + \mathbf{r}_0$ and \mathbf{r} denotes the total concentration of APC monomer. Here, by ignoring the ground state dynamics after pumping, \boldsymbol{e}_{gb} and \boldsymbol{e}'_{gb} can represent the absorption coefficients of GB for the states $|\mathbf{a}_{k}\rangle|\mathbf{b}_{k}\rangle$ and $|\mathbf{n}\rangle$, respectively. The terms \mathbf{e}_{ia} and \mathbf{e}'_{ia} are the absorption coefficients of IA for the states $|\mathbf{a}\rangle|\mathbf{b}_{s}\rangle$ and $|\mathbf{n}\rangle$, and \mathbf{e}_{se} and \mathbf{e}'_{se} are the absorption coefficients of SE for the states $|\mathbf{a}\rangle|\mathbf{b}\rangle$ and $|v\rangle$. Using the abovementioned equations $\Delta A = A - \overline{e}_{gb} \mathbf{r} = (\mathbf{e}'_{ia} - \mathbf{e}'_{se} - \overline{e}_{gb})\mathbf{r}_1 + (\mathbf{e}_{ia} - \mathbf{e}_{se} - \overline{e}_{gb})\mathbf{r}_2$ where $\vec{e}_{gb} = e'_{gb} + e_{gb}$. We can see that $\Delta \vec{A}$ provides information on the time-development of the populations of r_1 and r_2 via the GB, SE, and IA processes. We assume that the coefficients of \boldsymbol{e}_{gb} , \boldsymbol{e}'_{gb} , \boldsymbol{e}_{ia} , \boldsymbol{e}'_{se} and \boldsymbol{e}'_{se} are given by the Gaussian function, i.e. $\mathbf{e}_i = a_i \exp[-((\mathbf{w} \cdot \mathbf{w})/\Lambda_i)^2]$. Here *i* is the index representing each coefficient component, Λ_i is associated with the bandwidth, and a_i is the corresponding maximum intensity of each band. All the parameters are listed in table 1. Using these parameters, we successfully

simulate the three-dimensional transient absorption spectra as shown in figure 4 with a wavelength resolution of 5 nm and a time step of 1.5 ps. From figures 2a and 4, good agreement can be seen between the simulation spectra and the observed spectra in both time and frequency domains. In the simulation, we use two Gaussian functions centered at 587.5 nm and 618.7 nm for the absorption spectra of GB band, the IA band at the red side of the GB/SE band is described by the ground state absorption spectrum with a red shift of 1200 cm⁻¹, as shown in figure 5. In the calculated spectrum in figure 5b, we use a ratio of 1:1:0.83 for the GB, SE, and IA intensities, respectively. This IA spectrum is described by two Gaussian functions centered at 636.5 nm and 667.7 nm. Here, the dynamics of the IA band at the blue side of GB/SE band is still unknown, we temporally describe this IA band with one Gaussian function centered at 515 nm.

3.2b APC trimer: We find that our transient absorption spectra of APC trimer pumped at 580 nm agree with the reported in literature.⁸ An anisotropy profile reported in our previous work⁵ indicates that APC trimer should possess at least two electronically excited states in the visible region. The CD spectra, in fact, imply that the visible region of APC trimer involves at least three electronic states. However, the characteristic decay time of anisotropy profile is also found to be 200~500 fs upon excitations at 580 nm or 620 nm and probed at below 650 nm in this work. We could not obtain any anisotropy decay components when probed at shorter than 650 nm. This phenomenon provides us with an idea that with an excitation at 580 nm, hot vibrational states of the second electronic state are generated, from where fast internal conversion (IC) takes place within 200~500 fs.

For describing the three-dimensional, transient spectra of APC trimer, we need a model consisting of four states. In this model, subsequently to the excitation, the initial population is created in a hot vibrational state of the second excited electronic state $(|2_n\rangle)$ and via IC, the population relaxes from $|2_n\rangle$ to hot vibrational states of the first excited electronic state $(|1_{v'}\rangle)$ and $(|1_{v'}\rangle)$. The vibrational states $|1_{v'}\rangle$ and $|1_{v'}\rangle$ represent two vibrational modes with fast and slow relaxation rates respectively. Then the population relaxes

	SE position/width/intensity	IA position/width/intensity	GB position/width/intensity
n ⟩ vib. relax.	630·5 nm/18 nm/0·848 680·8 nm/48·2 nm/0·16	515·0 nm/35 nm/0·4 636·5 nm/45·7 nm/0·31 667·7 nm/23·4 nm/0·39	587·5 nm/45·7 nm/0·529 618·7 nm/23·4 nm/0·653
$ \boldsymbol{a}_e\rangle \boldsymbol{b}_g\rangle$ elec. relax.	635·5 nm/21 nm/0·848 680·8 nm/48·2 nm/0·31	515·0 nm/35 nm/0·4 636·5 nm/45·7 nm/0·31 667·7 nm/23·4 nm/0·39	

 Table 1a.
 The list of center positions, bandwidths and relative intensities of APC monomer.

Table 1b. The list of time constants (p	ps).	•
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$1/k_r$	$1/k_e$
25	300







Figure 5. The model of IA spectra of APC monomer. (a) The absorption and fluorescence spectra of APC monomer are assumed to result from an excitation of one subunit in the monomer. (b) Comparison between the calculated and observed TA spectra.

from the hot vibrational states to the bottom of the first excited electronic state (|1)). Since the fast IC dominates in the initial relaxation process, nonradiative transition (k_n) from a hot vibrational state to the bottom of the second excited electronic state (|2)) can be ignored. Three important rate constants $1/0.2 \text{ ps}^{-1}$, $1/2.5 \text{ ps}^{-1}$, and $1/30 \text{ ps}^{-1}$ are found in our fitting. In our model, we assign the rate constants as the IC rate k_{IC} , vibrational relaxation rates k_{rI} (fast) and k_{r2} (slower) of $|1_{v'}\rangle$ and $|1_{v'}\rangle$, respectively. The electronic lifetime $1/k_e$ of the first excited state $|1\rangle$ is assumed to be longer than 10 ns. For the trimer case, we have $d\mathbf{r}_{bb}/dt = -k_{IC}\mathbf{r}_{bb}$, $d\mathbf{r}_{cc}/dt = k_{IC}\mathbf{r}_{bb} - k_{r1}\mathbf{r}_{cc}$, $d\mathbf{r}_{dd}/dt = k_{IC}\mathbf{r}_{bb} - k_{r2}\mathbf{r}_{dd}$, and $d\mathbf{r}_{ee}/dt = k_{r1}\mathbf{r}_{cc} + k_{r2}\mathbf{r}_{dd} - k_e\mathbf{r}_{ee}$. Thus, the difference absorbance in transient spectra \mathbf{D} can be written as

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$$\Delta A = A - A_0 = (\boldsymbol{e}_{ia1} - \boldsymbol{e}_{se1} - \bar{\boldsymbol{e}}_{gb})\boldsymbol{r}_{bb},$$

+ $(\boldsymbol{e}_{ia2} - \boldsymbol{e}_{se2} - \bar{\boldsymbol{e}}_{gb})\boldsymbol{r}_{cc} + (\boldsymbol{e}_{ia3} - \boldsymbol{e}_{se3} - \bar{\boldsymbol{e}}_{gb})\boldsymbol{r}_{dd},$
+ $(\boldsymbol{e}_{ia4} - \boldsymbol{e}_{se4} - \bar{\boldsymbol{e}}_{gb})\boldsymbol{r}_{ee},$

where

$$\overline{\boldsymbol{e}}_{gb} = \boldsymbol{e}_{gb1} + \boldsymbol{e}_{gb2} + \boldsymbol{e}_{gb3} + \boldsymbol{e}_{gb4}$$

$$A_0 = \overline{\boldsymbol{e}}_{gb1} \, \boldsymbol{r}_{aa}(0) \, ,$$

and

$$\mathbf{r}_{aa}(0) = \mathbf{r}_{aa} + \mathbf{r}_{bb} + \mathbf{r}_{cc} + \mathbf{r}_{dd} + \mathbf{r}_{ee}.$$

Here, \mathbf{r}_{aa} , \mathbf{r}_{bb} , \mathbf{r}_{cc} , \mathbf{r}_{dd} and \mathbf{r}_{ee} represent the populations of ground state, $|2_{\mathbf{n}}\rangle$, $|1_{v'}\rangle$, $|1_{v'}\rangle$ and bottom of the first excited state. Figure 6 shows the calculated three-dimensional, TA spectra of APC trimer with a wavelength resolution of 3 nm and a time step of 0.3 ps, which can be compared with figure 2b. In the simulation, the GB band is described by three Gaussian functions centered at 652.3 nm, 634.5 nm, and 602.7 nm; the IA band is described by two Gaussian functions at 647 nm and 680 nm. All the parameters are listed in table 2. It is interesting to note here that the 680 nm IA band results only from the second excited state $|2_{\mathbf{n}}\rangle$ and the IA band at 647 nm appears at first excited state. The peak position of the SE band of the $|1_{v'}\rangle$ state exhibits a slightly red

 Table 2a.
 The list of center positions, bandwidths and relative intensities of APC trimer.

	SE position/width/intensity	IA position/width/intensity	GB position/width/intensity
2 n >	730 nm/36 nm/0·13	680 nm/30 nm/0·15	602·7 nm/44·8 nm/0·12 634·5 nm/30 nm/0·11 652·3 nm/11·6 nm/0·19
$ 1_{n'}\rangle$	658 nm/10 nm/1·25 675 nm/10 nm/0·25 730 nm/30 nm/0·187	647 nm/20 nm/0·05	
1 " ″⟩	659 nm/12 nm/1·5 679 nm/10 nm/0·2 730 nm/30 nm/0·19	647 nm/20 nm/0·3	
1>	660 nm/12 nm/1·13 679 nm/13 nm/0·37 730 nm/30 nm/0·1	647 nm/20 nm/0·44	

1/k _{IC}	$1/k_{rl}$	$1/k_{r2}$	$1/k_e$
0.2	2.5	30	10000



observed spectra at various delay times: 0 ps, 1 ps, 5 ps, and 30 ps. (b) Comparison between calculated and observed dynamic profiles at various wavelengths: 635 nm, 660nm, 670nm, and 680nm. a) Comparison between calculated and Figure 6. Calculated three -dimensional, transient absorption spectra of APC trimer. (

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shift to 679 nm. Figures 6 show a comparison between the simulation and observed spectra in both time and frequency domains and good agreement can be seen. We find that the fast dynamics of APC trimer results from IC process. Basically, six exciton states can exist once the electronic interactions among six chromophores are considered. However, through the fast non-radiative channel IC, no matter which state is excited the population will relax to the first excited state during 200~500 fs. The first excited state is the only one electronic state which can trap the final population.

4. Conclusion

We have shown our transient absorption spectra of APC. To understand the observed TA spectra, we have constructed a model for both the energy levels and the dynamics. Based on the proposed model, for the monomer case, we have obtained an IA band with a red shift of 1200 cm⁻¹ from the ground state absorption band. This indicates that the energy gap between the ground and excited states of the unexcited subunit decreases by 1200 cm⁻¹. For APC trimer, we have found that the excited population quickly arrives at the first excited state via IC, which allows us to observe the very fast dynamics within the early pumping time. In our model, we have found that for APC monomer, we need at least two excited states to explain the all processes while for APC trimer, at least four excited states are needed. The purpose of constructing model is not to reproduce the TA results accurately but to obtain the essence of dynamics. Our models can provide essential pictures for understanding the transient dynamics of APC trimer and monomers.

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